

Physics of Scale Selection for Mesoscale Patterns in Drift Wave Turbulence

Special Focus: Physics of ExB Staircase (楼梯)

Weixin Guo¹, P. H. Diamond², Arash Ashourvan³

1. IFPP Lab, SEEE, HUST
2. CASS and Dept. of Physics, UCSD; SWIP
3. PPPL; DIII-D

***Supported by U. S. Department of Energy, Office of Science, Office of Fusion Energy Sciences, under Award Number DE-FG02-04ER54738**

***Supported by Graduate School, HUST and NSFC Grant Nos. 11675059 and 11305071**

■ Basics of nonlinear patterns

- Pattern formation (**scale, structure ?**) in magnetic confinement plasma ↔ **Scale selection**

Especially, zonal flow → staircase

- Formation mechanism and **feedback loops**

→ **Shearing** and **Rhines** mechanisms, transport bifurcation,

■ Recent model study of zonal pattern formation / sustainment

- ✓ Parameter scans
- ✓ Feedback loop studies
- ✓ Turbulence spreading
- ✓ Initial condition sensitivity
- ✓ Boundary condition sensitivity.....

■ Conclusions and plan

Scales Determine Transport



■ Several spatial scales enter transport

- ✓ Drift wave, micro-scale ρ
 - ✓ Mean flow, macro-scale $L_n \sim a$
 - ✓ Zonal flow, meso-scale ($\sqrt{\rho L_n}$?)
- ρ : Larmor radius
 - a : minor radius
 - $L_n = -n/\nabla n$: density scale length

■ Goal→Predicting turbulence and transport in saturated states

$$D = D_{Bohm} \left(\frac{\rho}{a} \right)^\alpha, \quad 0 < \alpha < 1. \quad \begin{cases} \alpha = 0, \text{ Bohm scaling (Bad)} \\ \alpha = 1, \text{ Gyro-Bohm scaling (Good)} \end{cases}$$

→ Turbulent diffusivity D scaling with ρ/a is a key question in fusion!

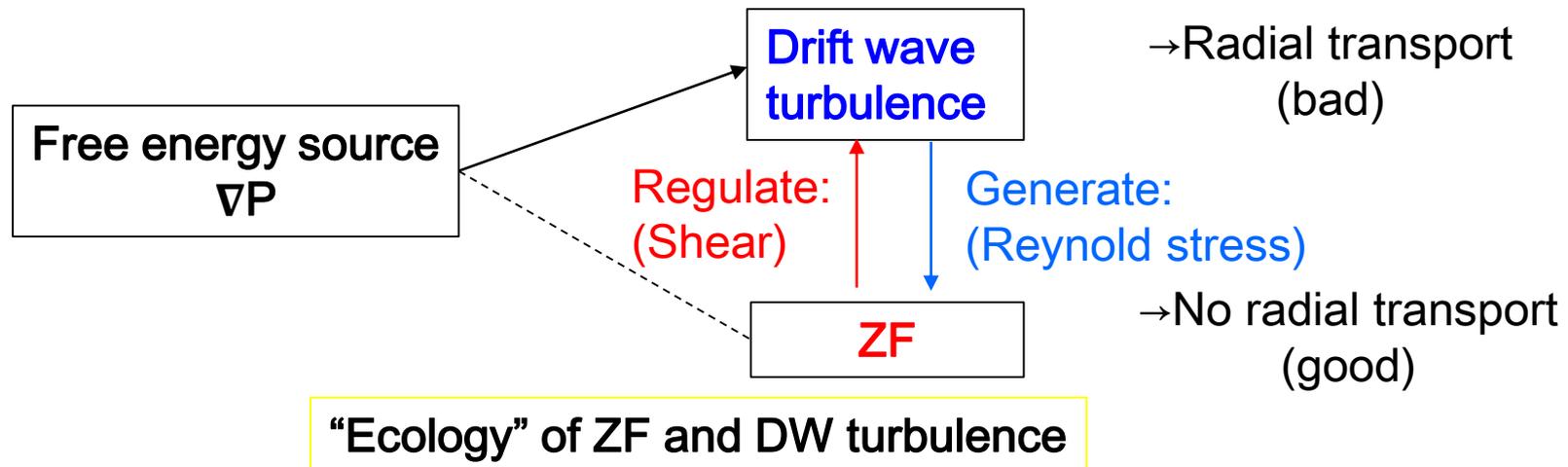
→ Value of α is directly linked to **meso-scale selection**

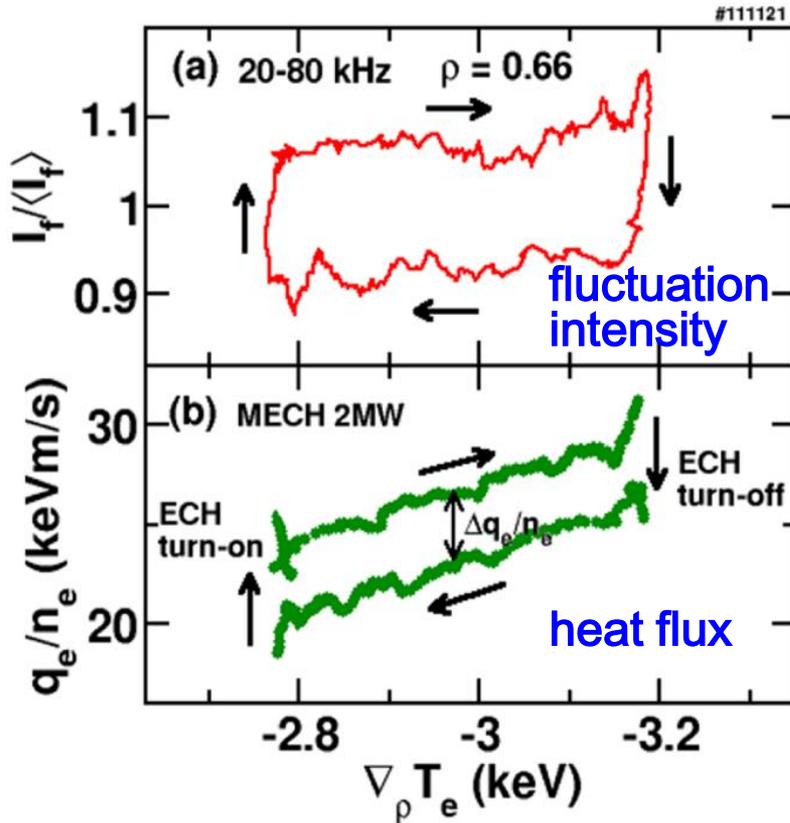
■ The ecology of feedback loops must enter the scaling of **spatial structure**

$$|\delta\varphi| \sim l_{mix} \sim \frac{1}{1 + \sigma(v'_{E \times B})^2}$$

Physics of Scale Selection

- **Formation of patterns** in self-organizing, non-equilibrium and nonlinear systems is widely observed.
 - ✓ Layered Stratification
 - ✓ Potential Vorticity (PV) and ExB Staircase
 - ✓ Transport Barrier
 - ✓ Zonal Flow (ZF)
- ZF:
 - minimal inertia and Landau damping → benign repository for free energy
 - **regulate** turbulent transport, trigger L-H and ITBs
- **Closing the feedback loop** when predators meet the prey





[S. Inagaki, NF 2013]

LHD, L-mode

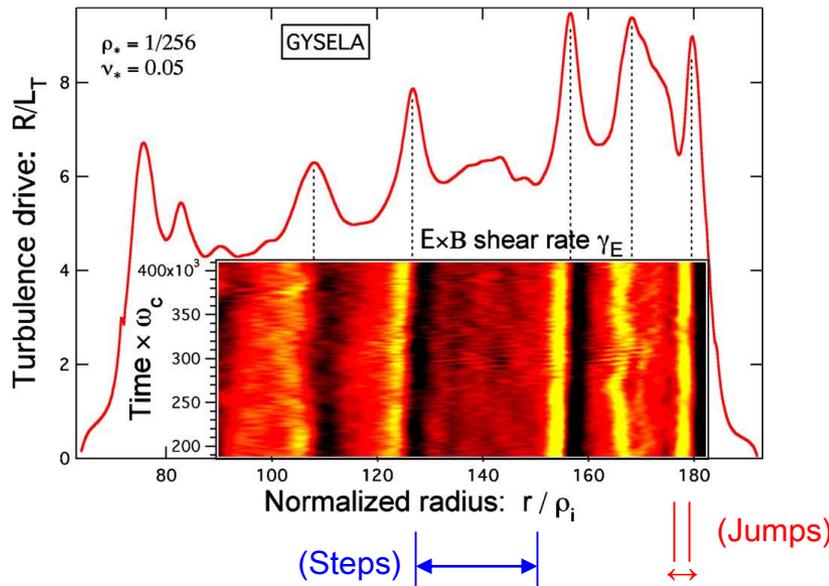
- Experimental observation of **hysteresis** reflected by the non-linear ∇T_e dependence of fluctuation intensity (a) and heat flux (b).
 - self-organized **global hysteresis**
- Hysteresis strongly suggests **bistability** in the fluctuation intensity
- Here, not due to local transport barrier

A Hint: ExB Staircase

- ✓ In ion temperature gradient (ITG) turbulence, **quasi-regular (spatial) and long-lived (temporal) E×B flows coexisting with temperature corrugations** are observed numerically + experimentally

→ Meso-scale **E×B staircase (楼梯)**, by analogy to PV ($\overset{\text{density}}{\uparrow} n - \overset{\text{vorticity}}{\uparrow} u$) staircases and atmospheric jets

Dritschel, McIntyre, Journal of the Atmospheric Sciences, 2008



“A natural and dynamic means for the simultaneous existence of these **two antagonistic trends**”

- ✓ Clear scale selection
- ✓ Clear link between **ZF scale** ↔ **avalanche scale**

$\delta^{\text{flow}} \sim 10 \rho_s$

(Jumps)

$\Delta^{\text{stat}} \sim 40 \rho_s$

(Steps)

[G. Dif-pradalier, *et al*, Physical Rev. E **82**, 025401 (2010)]
 [G. Dif-pradalier, *et al*, , Phys. Rev. Lett. **114**, 085004 (2015).]

- Works: simulations in ITG turbulence (GYSELA, GKNET)

G. Hornung et al NF 2017

G. Dif-Pradalier et al NF 2017

W. Wang, Y. Kishimoto et al NF 2018

Leave many questions:

→ Conditions for existence of the $E \times B$ staircase?

→ Explore **dimensionless parameter** dependence?

→ Vary **boundary** condition (B. C.) and **initial** condition?

→ **Robustness** in parameter space?

→ **Impact** on transport and confinement?

→ **Relation** to the barrier?

- Especially, **which feedback is critical** for the formation and sustainment of $E \times B$ staircase?

Reduced models: Why and What ?



- **Reduced models:** self-consistently relate variations in **mean plasma fields** to **fluctuation intensity** !

lower computational cost, flexibility, simplicity, variety, supply **essential** understanding.....

- Hasegawa–Wakatani (H–W) system

✓ Simple **generic** system describing collisional drift wave turbulence, which conserves energy and **P**otential **E**nstrophy \leftrightarrow $PE \equiv \frac{1}{2} (n - u)^2$

density vorticity

✓ **Inhomogeneous PV mixing and symmetry** generate ZF from turbulence via transport

- Flexible reduced model to understand the physics of meso-scale patterns derived from H-W equation

✓ Inhomogeneous effect of the mean field (density, PV.....)

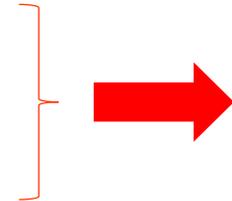
✓ **Mixing length: key element** → What is it ?

Mixing Length → Rhines Scale



- Mixing length is a nonlinear hybrid of two length scales

- A constant **excitation** scale l_0
- A **dynamic** length scale l_d
a function of the **system gradient**


$$l_{mix}^{\kappa} = \frac{l_0^2}{1 + l_0^2 / l_d^2}$$

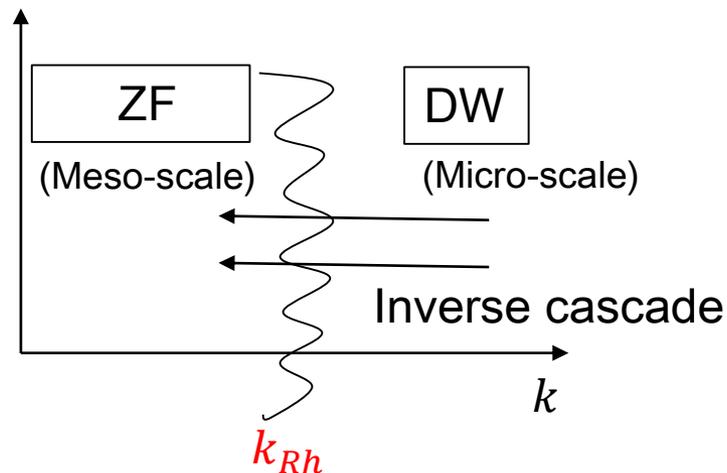
- Not the conventional mode scale B B Kadomtsev 《Plasma turbulence》, 1965
- Inhomogeneous PV mixing process selects the mixing scale l_{mix} to generate pattern.
→ Feedback loop strongly depends on l_{mix}
- $l_d = l_{Rh}$ ↔ Rhines scale

What is Rhines Scale ?

- Scale where **decorrelation rate** crosses **mismatch frequency** (due to dissipation)

(wave coupling mismatch)

$$k_{Rh} \tilde{v} = \omega_{MM} \quad \rightarrow \quad l_d = l_{Rh} = \frac{\sqrt{\varepsilon}}{|\partial_x(n-u)|}$$



- ε : fluctuation PE
- $\partial_x(n-u)$: mean PV gradient

$$l_{mix} = \frac{l_0}{\left(1 + l_0^2 \left[\partial_x(n-u)\right]^2 / \varepsilon\right)^{\kappa/2}}$$

- $l < l_{Rh} \rightarrow$ short memory, turbulence is eddy-like \rightarrow strong mixing
- $l > l_{Rh} \rightarrow$ long memory, turbulence is wave-like \rightarrow weak mixing

Reduced Model from H-W Equation



- Evolution of mean density n , vorticity u and turbulent potential enstrophy ε

$$\partial_t n = \partial_x D_n \partial_x n + D_c \partial_x^2 n, \quad (1)$$

$$\partial_t u = \partial_x (D_n - \chi) \partial_x n + \chi \partial_x^2 u + \mu_c \partial_x^2 u, \quad (2)$$

Adiabatic limit

$$\partial_t \varepsilon = \partial_x D_\varepsilon \partial_x \varepsilon + \chi [\partial_x (n - u)]^2 - \varepsilon_c^{-1} \varepsilon^{3/2} + \mathcal{P}. \quad (3)$$

A. Ashourvan and P. H. Diamond, PoP 2017, PRE 2016

- $D_n \approx l^2 \frac{\varepsilon}{\alpha}$
 $\chi(x) = c_\chi l^2 \frac{\varepsilon}{\alpha}$
 $D_\varepsilon(x) \cong \beta l^2 \varepsilon^{1/2}$

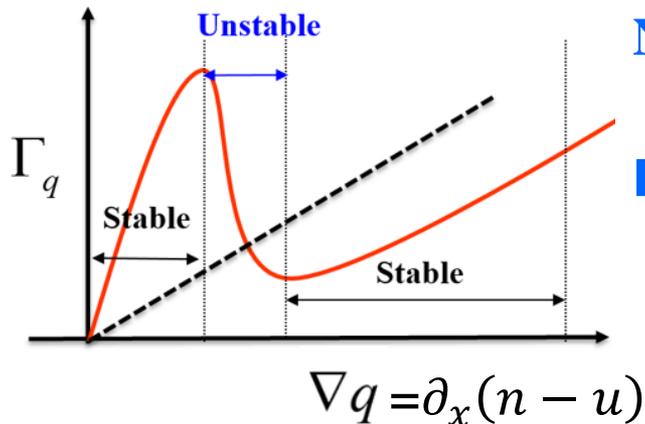
$$l_{mix} = \frac{l_0}{\left(1 + l_0^2 \left[\partial_x (n - u) \right]^2 / \varepsilon \right)^{1/2}}$$

- $PV = q = n - u$

Steepen of mean PV gradient

Further drop of l_{mix}

Drop in local ε

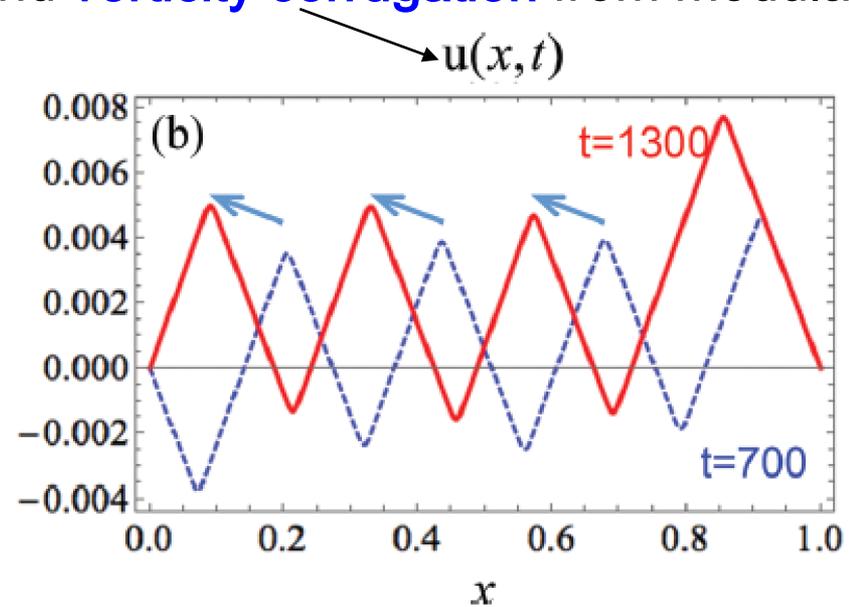
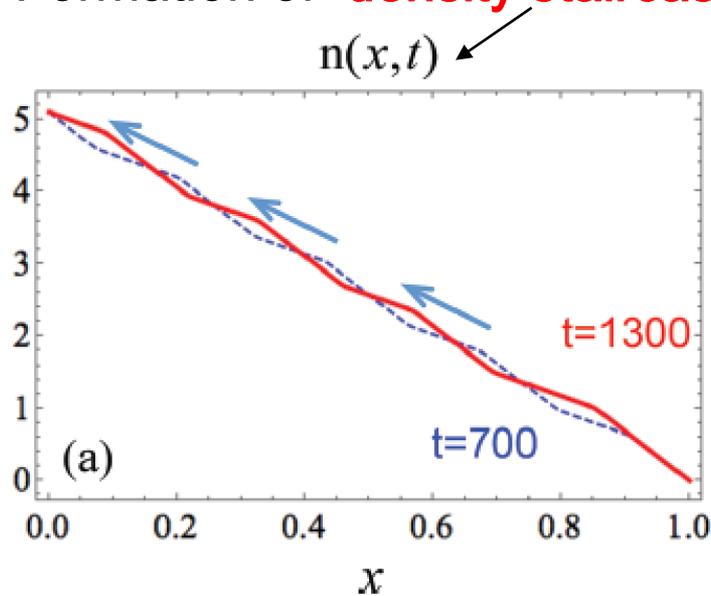


Negative diffusion region---bistable mixing

- ✓ The positive feedback loop
 - Drives the pattern
 - **Leads to nonlinear feature** formation in the mean profile.

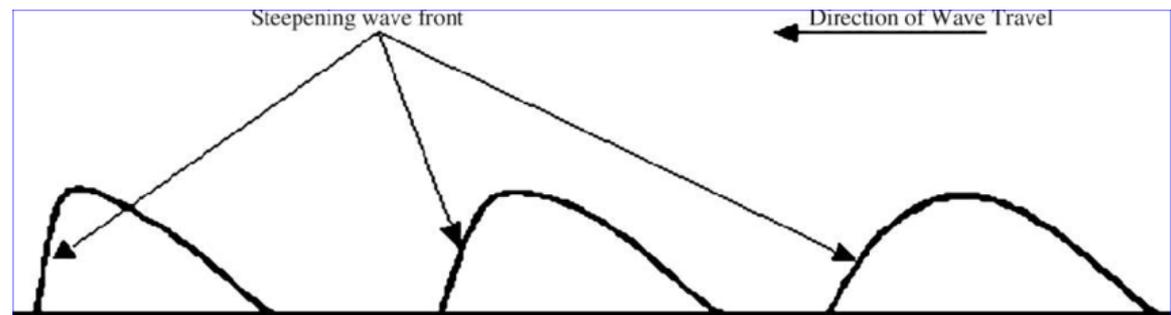
Reduced Model → Pattern Formation

- Formation of **density staircase** and **vorticity corrugation** from modulation



→ **“Staircase”** pattern : consequence of modulational bistability feedback

- Similar to shock wave, which forms due to **self-steepening** of ordinary waves



- What are the **dimensionless parameters** and how does the pattern structure respond to parameter scan?
- What is the **principal feedback loop** physics?
- How does the pattern respond to **turbulence spreading** and/or avalanching?
- Does the pattern have **memory** of initial condition?
- What is the effect of **mean $E \times B$ shear** on pattern?
(Macro \leftrightarrow Meso scale)

.....

Parameter sensitivity studies



■ Dimensionless parameters

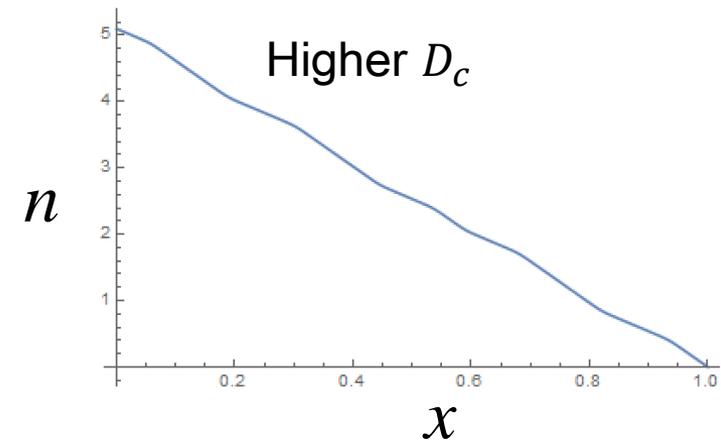
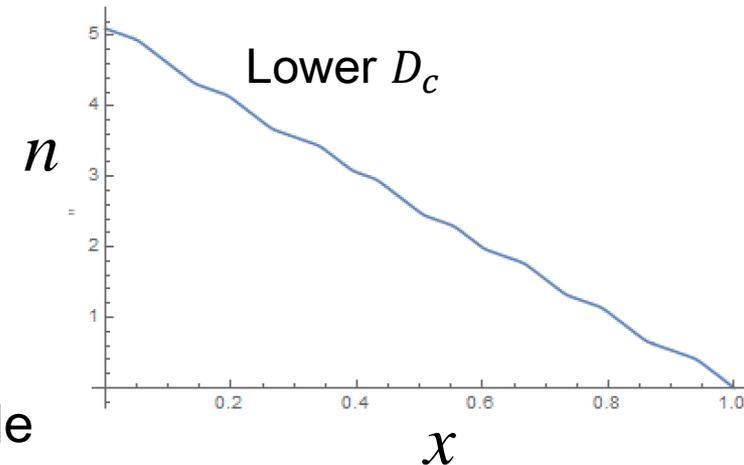
- $Re \sim \frac{\sqrt{\epsilon} l^2}{\mu_c}$
- $\alpha = \frac{k_{\parallel}^2 v_{th}^2}{\omega v_e} > 1$, fixed
- $Pr \sim \frac{\mu_c}{D_c}$
 - Collisional viscosity (pointing to μ_c)
 - Collisional diffusivity (pointing to D_c)
- $\frac{Drive}{dissipation} \sim \frac{\epsilon_c \gamma_{\epsilon}}{\sqrt{\epsilon}}$

■ Re and Pr scans:

FOM : **Number of steps** (N_s) in density profile

↔ structure complexity

- ✓ Both $\mu_c \nearrow$ and $D_c \nearrow$ (remove energy)
 - ➔ $N_s \searrow$ ➔ Damp the staircase
- ✓ $Pr \searrow$ ➔ $N_s \searrow$
- ✓ $\gamma_{\epsilon} \nearrow$ ➔ $N_s \nearrow$



What is the Feedback Mechanism?

- In the case of mean perpendicular shears, mixing length is presented

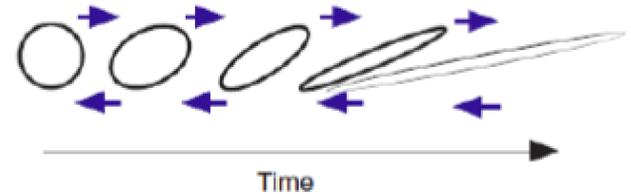
$$l_{mix}^2 = \frac{l_0^2}{\left[1 + (\bar{v}'_{E \times B})^2 \tau_c^2\right]^\kappa}$$

- l_0 : the mixing scale without $v_{E \times B}$
- τ_c : fluctuation correlation time
- $\bar{v}'_{E \times B}$: perpendicular shear rate

- The correlation time is given

$$\tau_c = (u^2 \varepsilon / l_0^2)^{-1/4}$$

- ε : turbulent PE
- $u = \bar{v}'_{E \times B}$



→
$$l_{mix} = \frac{l_0}{\left[1 + \frac{|u|}{l_0 \sqrt{\varepsilon}}\right]^{\kappa/2}}$$

↔ Feedback loop works through **shearing dependent** mixing length

- Rhines scale:
$$l_{mix} = \frac{l_0}{\left(1 + l_0^2 \left[\frac{\partial_x (n-u)}{\varepsilon}\right]^2\right)^{\kappa/2}}$$

∇n ← $\frac{\partial_x (n-u)}{\varepsilon}$ → $\partial_x u$

A. Ashourvan and P. H. Diamond,
PoP 2017, PRE 2016

- Multiple candidate, which mechanism is key?

Outcome (parameters as before)



- If only $E \times B$ shearing feedback ($u = \bar{v}'_{E \times B}$)

→ **No** layer staircase structure forms !

- If Rhines scale feedback (∇n and $\partial_x u$)

→ **Recover** staircase pattern, three stages:

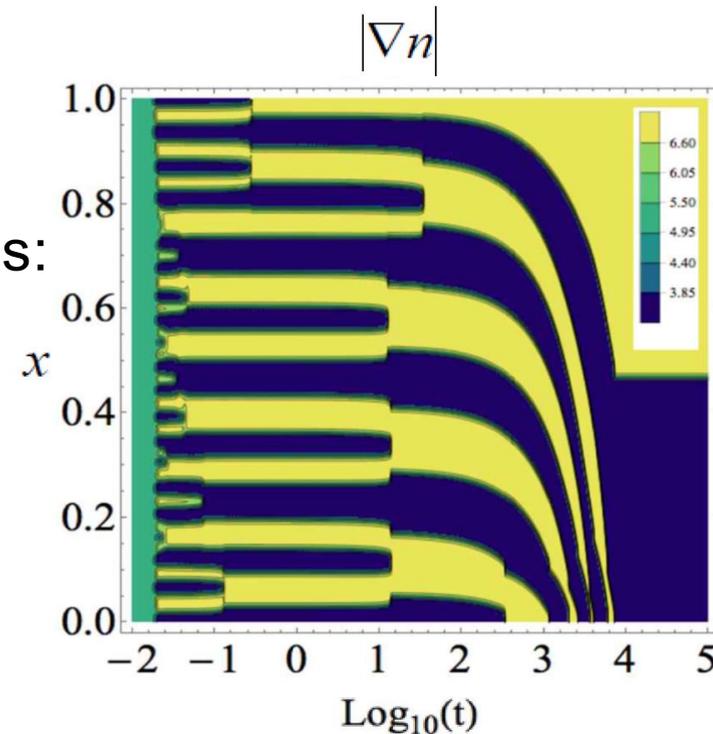
- Microscale instabilities → Non-linear (NL) mesoscale structure
- → Merger
- → Migration

→ Check $\partial_x u$!

- If turn off $\partial_x u$ → **Recover** pattern

↔ Experimental, Jiang Min

➔ Feedback through NL **gradient drive** of mixing **is key loop** !



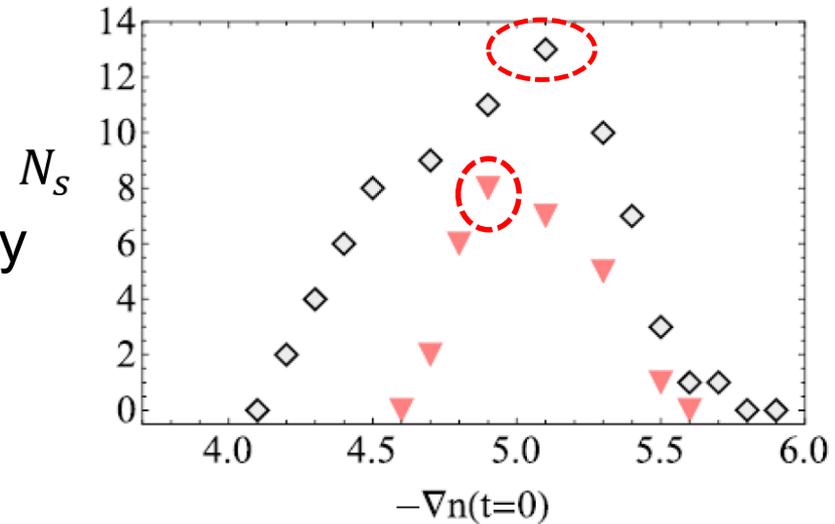
Sensitivity to (Initial) Gradient Drive



■ Scan initial gradient drive

$$\nabla n(t=0)$$

- Initial rise \leftrightarrow increased free energy to staircase structure $\rightarrow N_s \uparrow$
- Further rise $\rightarrow N_s \downarrow \leftrightarrow$ effects of diffusive dissipation limiting small scales

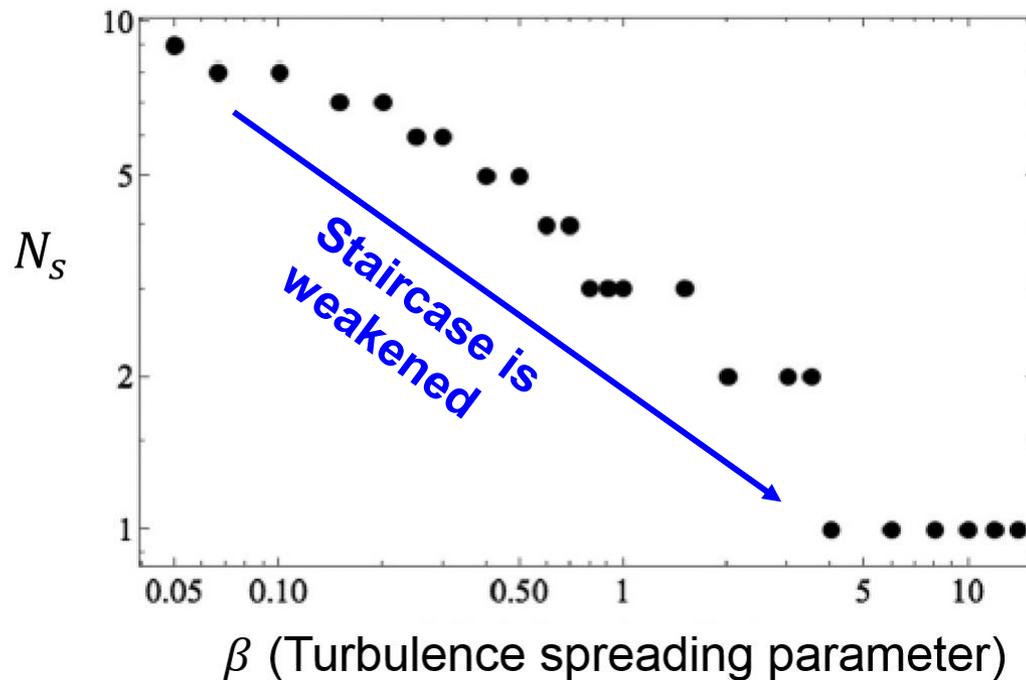


- ➔
- ✓ Free energy and dissipation define pattern structure
 - ✓ **Minimal step scale** exists

■ Flux Drive? What defines the selected scale? (Ongoing)

■ Evolution of fluctuation potential enstrophy (PE)

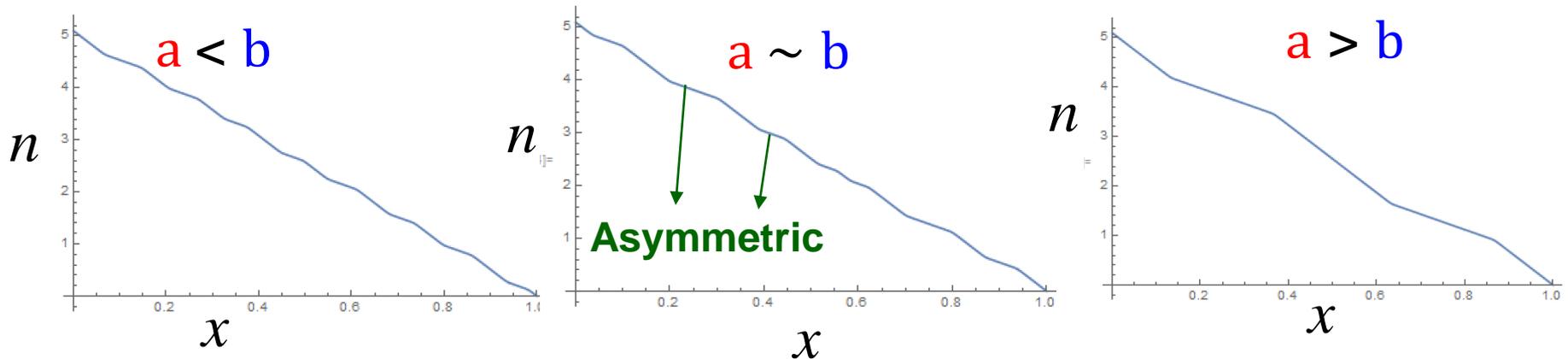
- β measures the effect of turbulence spreading of **PE ~ fluctuation intensity**
- β represents complex mode interaction physics



- ✓ Moderate increase in β **weakens** the staircase
- ➔
- ✓ Pattern is **sensitive** to turbulence spreading

- Initialize: $u(t = 0) = \nabla_{\perp}^2 \phi = a \sin(n\pi x) + b$
 - $a \sin(n\pi x)$: initial zonal shear
 - b : mean shear

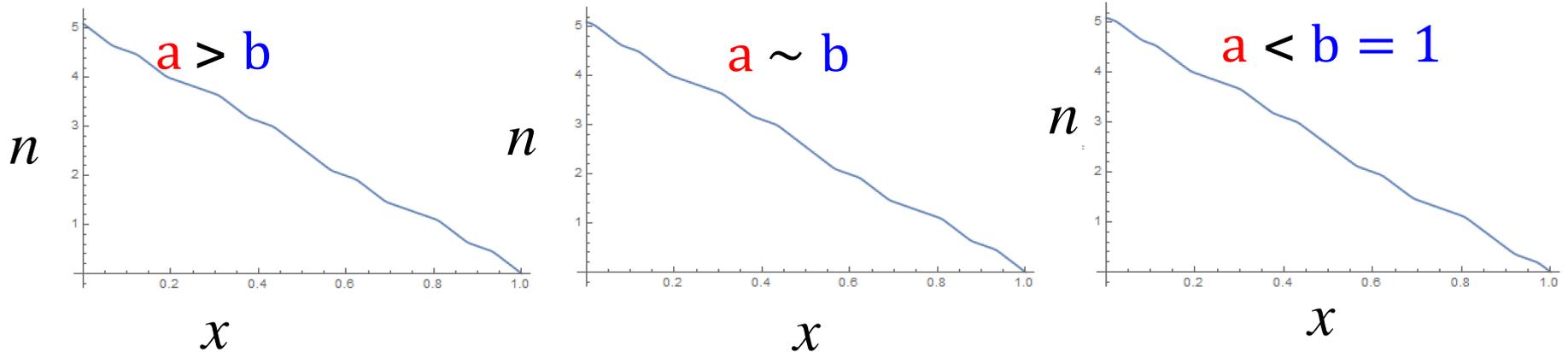
I. Fixed mean shear b , vary zonal shear a



- ✓ Asymmetric pattern can form when zonal shear is comparable with mean shear
- ✓ Stronger zonal shear ($a > b$) will destroy pattern

- Initialize: $u(t = 0) = \nabla_{\perp}^2 \varphi = a \sin(n\pi x) + b$
 - $a \sin(n\pi x)$: initial zonal shear
 - b : mean shear

II. Fixed zonal shear a , vary mean shear b



- ✓ $b \leq 1$, increasing mean shear leaves pattern unchanged
- ✓ Stronger mean shear ($b > 1$) destroys the pattern

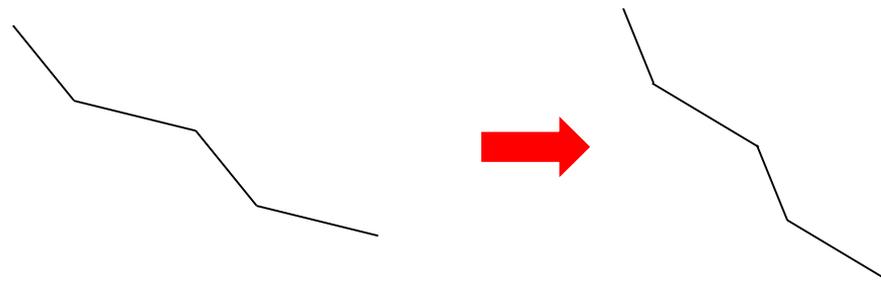


Pattern retains memory of initialization. Both zonal and mean shear are relevant !

In shearing feedback loop:

- Lower the damping value adopted in Rhines feedback loop
 - Recover the pattern
 - But, sustainment of the pattern is difficult
- Increase (a) the $\nabla n(t = 0)$ and (b) the initial period of zonal shear
 - N_s : first increase and then decrease !
- Increasing the zonal shear causes the jumps to become much

steeper !



Summary of results



Existence of density staircase and vorticity corrugation in H-W system

- Pattern responds to the **plasma parameters**
 - Increase flow **viscosity** (μ_c), particle diffusion **damping** (D_c):
 - **damp** pattern.
 - Moreover, D_c is more effective
 - Increase production → selects minimal step scale
- **Feedback loop physics**

Shearing vs Rhines → NL ∇n dependence of l_{mix} is key loop!

Surprise ! : shearing feedback is **ineffective**
- Moderate turbulence spreading **weakens** the Pattern
- Pattern **retains** the memory of initial scale
- Both **mean** and zonal $E \times B$ shear act on zonal pattern

- Origin of staircase is simple: quasi-periodic zonal pattern in u , n formed by **self-sharpening of modulation**
- ❑ Principal feedback is through gradient nonlinearity of mixing and transport (**here, ∇n determines**)
- ❑ Pattern **scale** emerges as **minimum** set by ∇n , dissipation, period of zonal shear.....
- Reduced models are **useful** complements to large scale simulation !

◆ Understanding

- Why shearing feedback is not effective?
- Requirement of driving gradient nonlinearity to sustain pattern?
- Other feedback mechanism?

◆ Flux driven studies (Extend Ashourvan, Diamond studies)

- Scaling of step size ?
- Quantity the global hysteresis of state ?
- Determined condition for ITB transition ?
(global→local barrier ?)

◆ Modify the model to treat mean shear consistently, and to treat high density (i.e., hydrodynamic) regime ?.....

**Thank you very much for
your attention!**

**International Joint Research Laboratory of Magnetic
Confinement Fusion and Plasma Physics
(IFPP, J-TEXT)**

Wuhan, China