

The ubiquitous zonal flows (and corrugations)

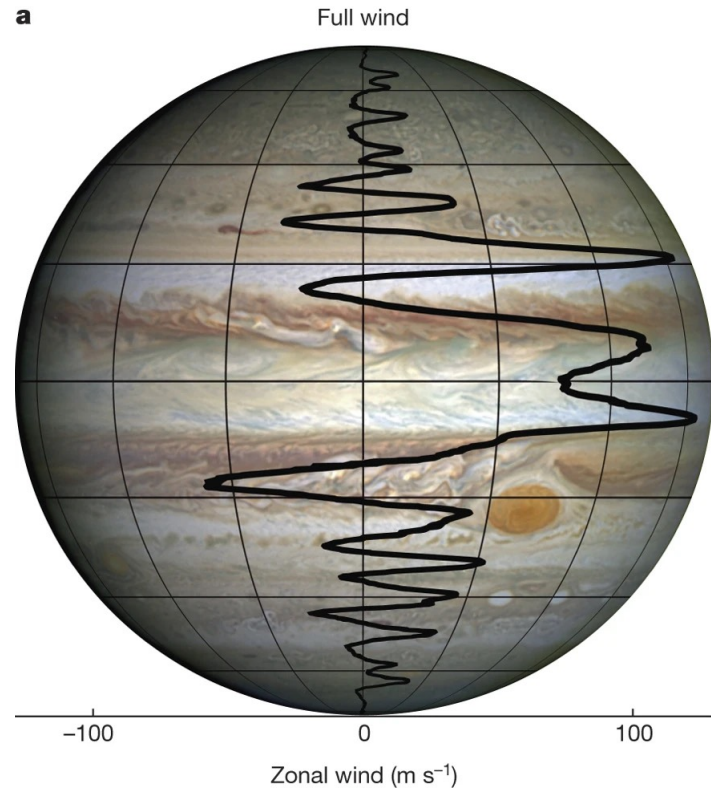
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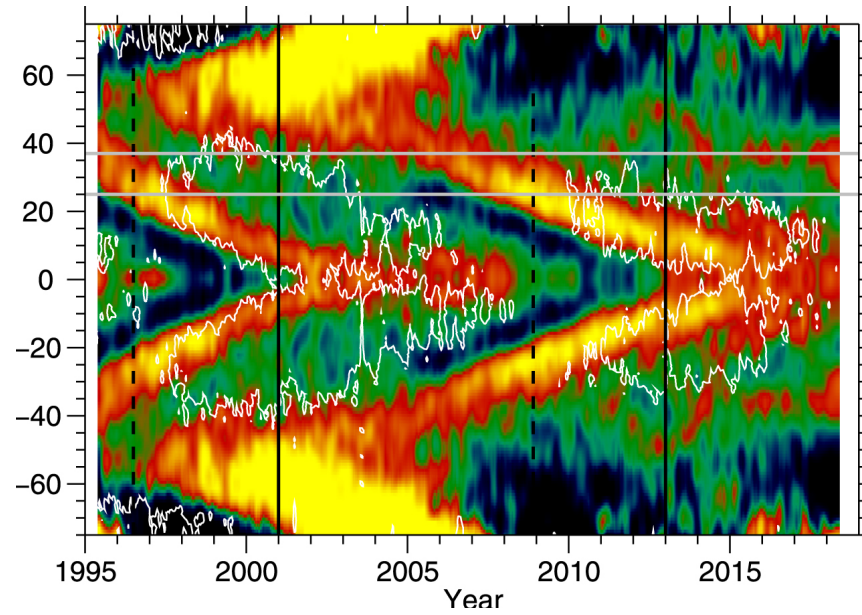
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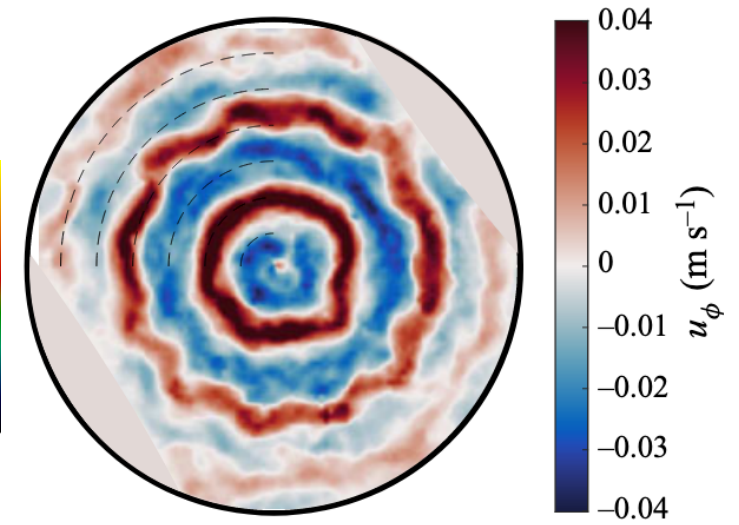
Zonal flows are ubiquitous



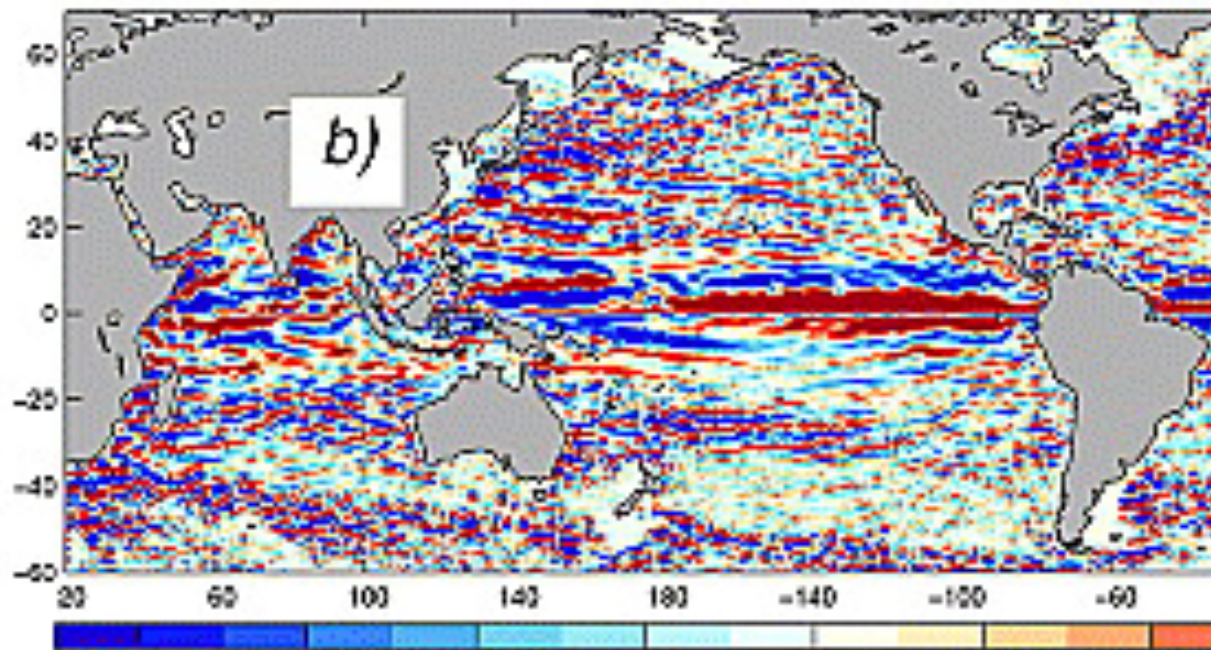
ZFs in Jupiter



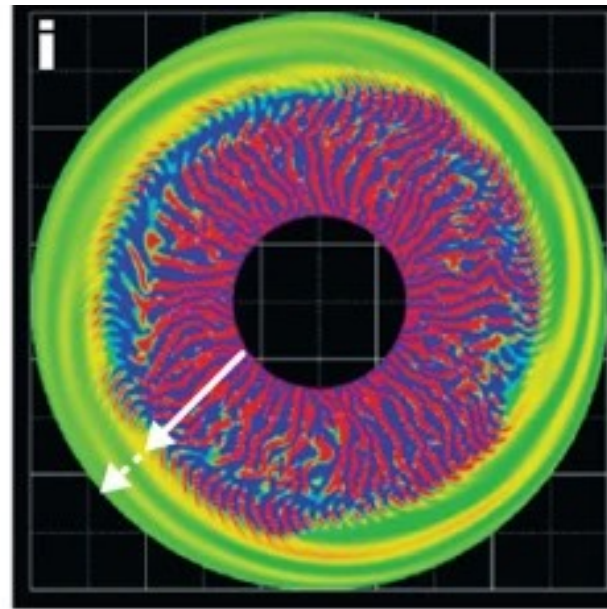
Migrating ZFs in convection zone of Sun



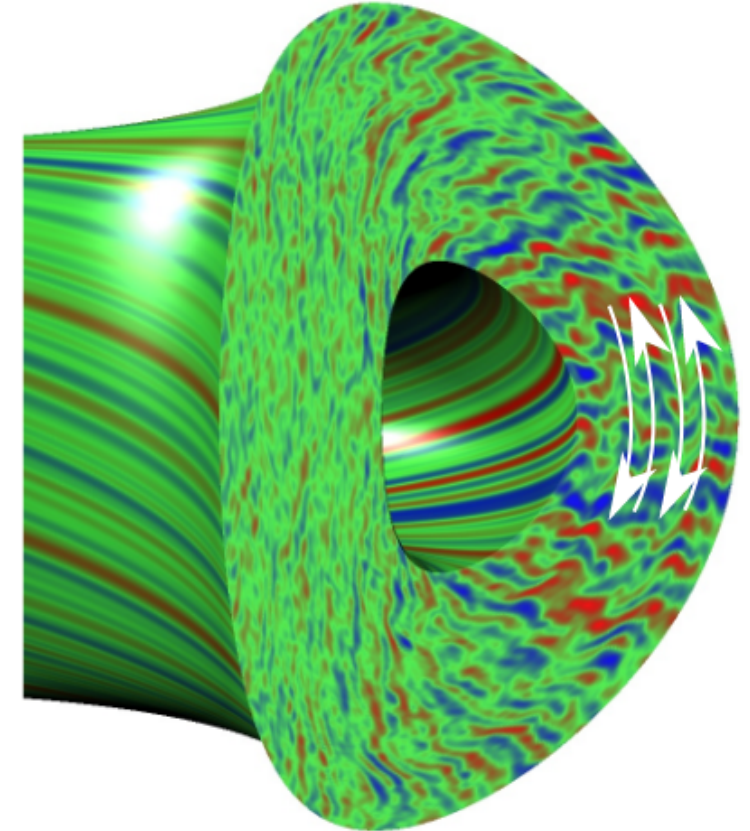
ZFs in rotating liquid column experiment
Lemasquerier *et al* JFM



ZFs in ocean Maximenko *et al* GRL 2005



ZFs in liquid interior of Earth outer core T Miyagoshi *et al* Nature 2010



ZFs in tokamak

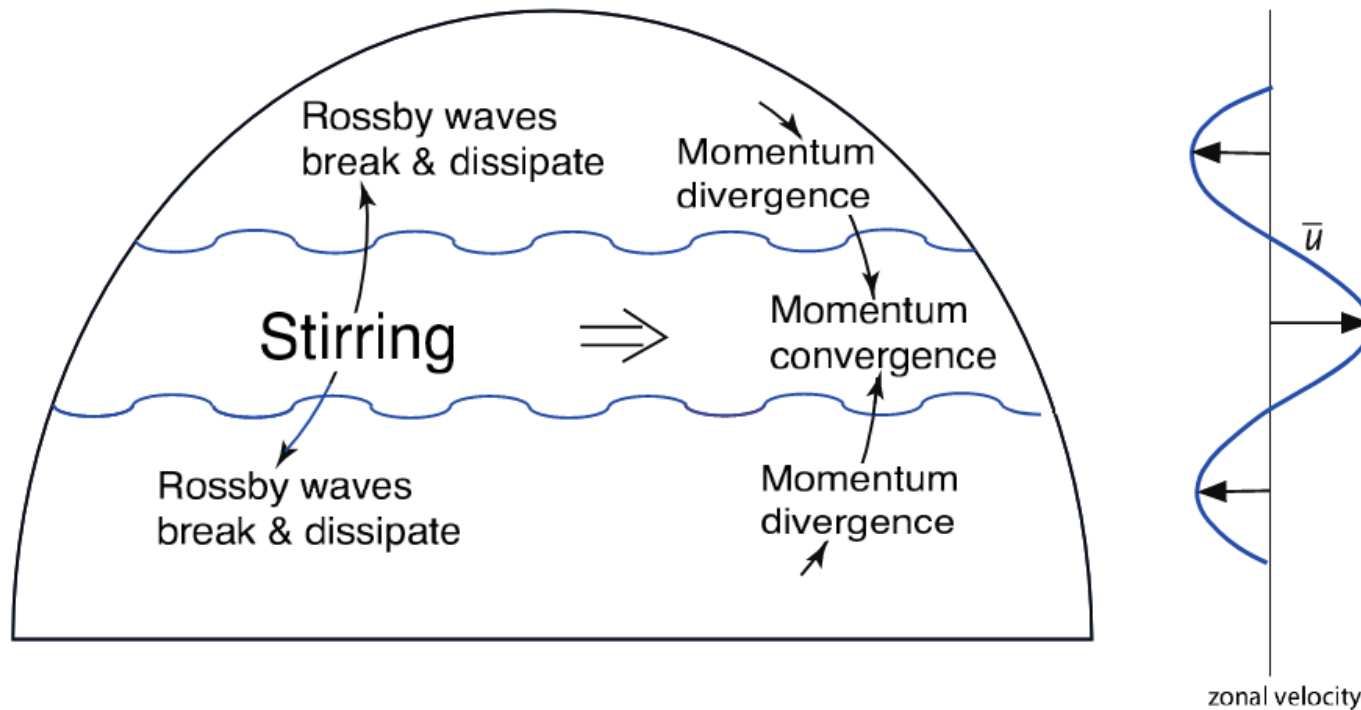
Outline

- Why zonal flows are ubiquitous?
- What determines the ZF pattern: Momentum theorem
- Shearing effects
 - Turbulence decorrelation, Induced diffusion ,
Modulational instability
- Unifying zonal flows and corrugations
 - Noise + modulations
 - Zonal cross correlation- staircase
- Feedback loops (noise effects)
- Conclusions

Why zonal flows are ubiquitous?

GFD Perspective

- Mid latitude zonal circulation [G K Vallis Book]



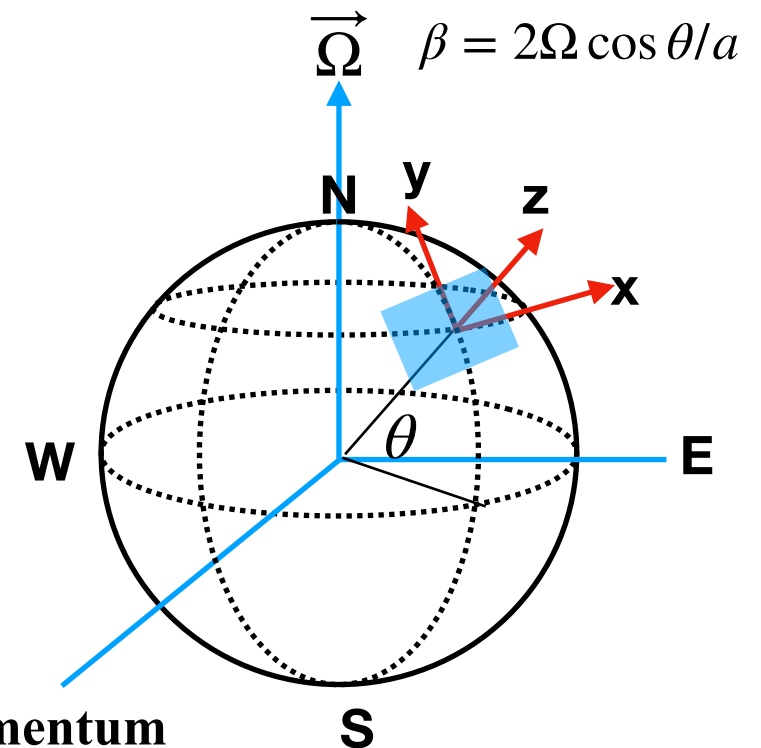
- Stirring in mid-latitudes (by baroclinic eddies) generates Rossby waves that propagate away.
- Momentum converges in the region of stirring, producing eastward flow there and weaker westward flow on its flanks.

Rossby waves on beta-plane

- Frequency $\omega = -\frac{\beta k_x}{k_x^2 + k_y^2}$, Group velocity $v_{gy,k} = \frac{2\beta k_x k_y}{(k_x^2 + k_y^2)^2}$

- Momentum flux: $\langle v_y v_x \rangle = -\sum_{\vec{k}} \frac{1}{2} k_x k_y |\psi|^2 = -\sum_{\vec{k}} v_{gy,k} \mathcal{P}_k$

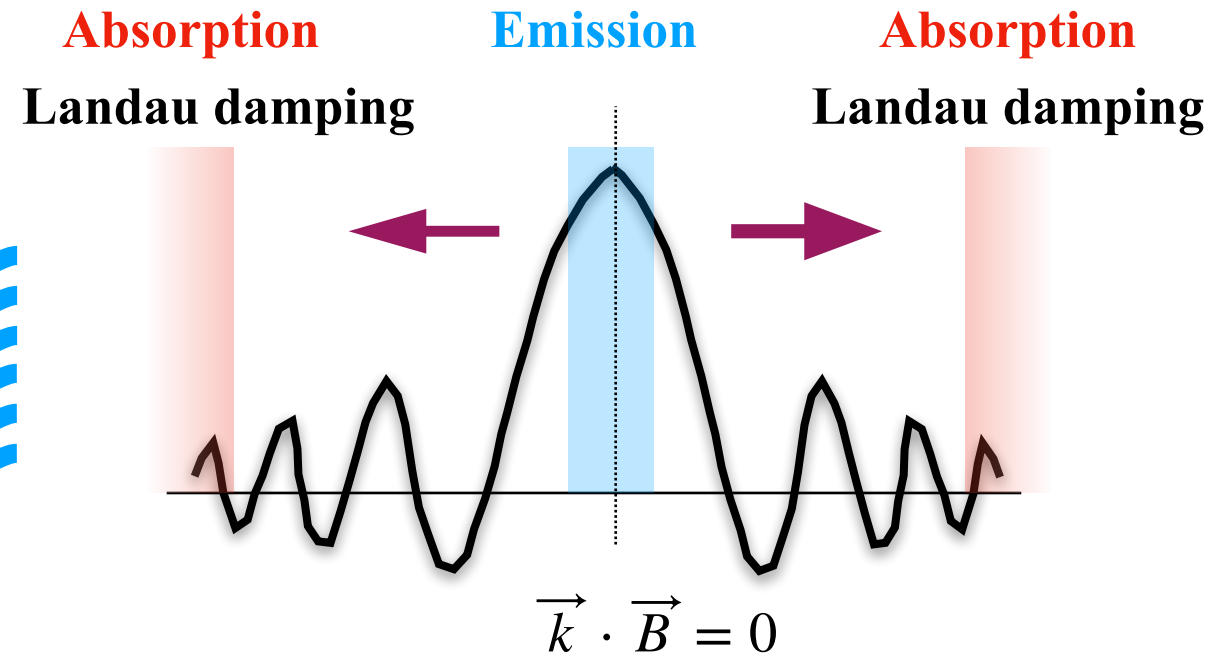
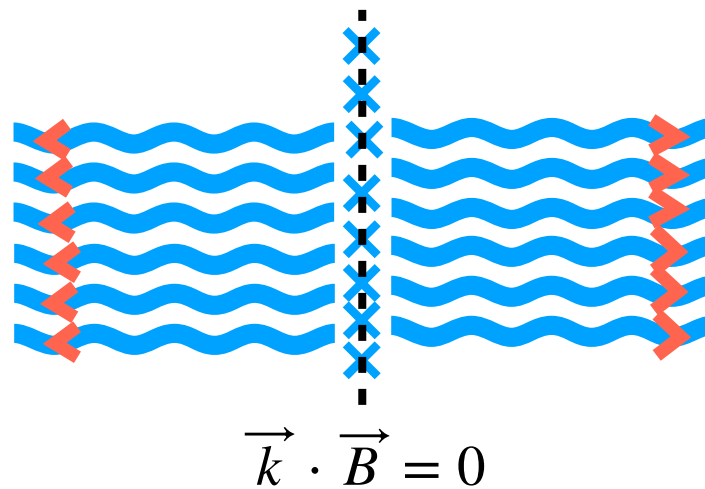
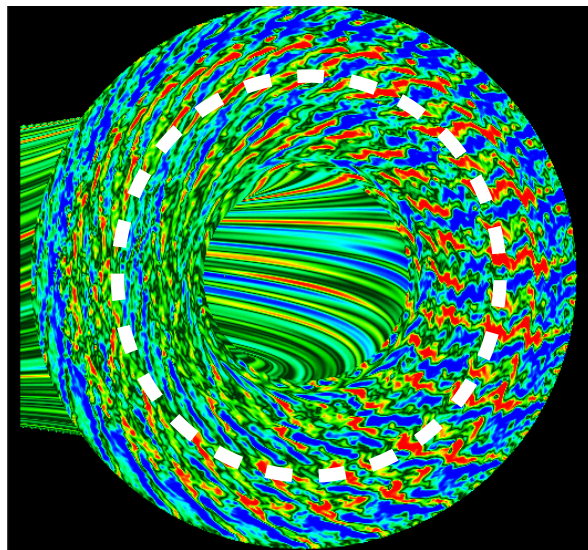
• Outgoing waves \implies Incoming momentum flux



Why zonal flows are ubiquitous?

MFE perspective

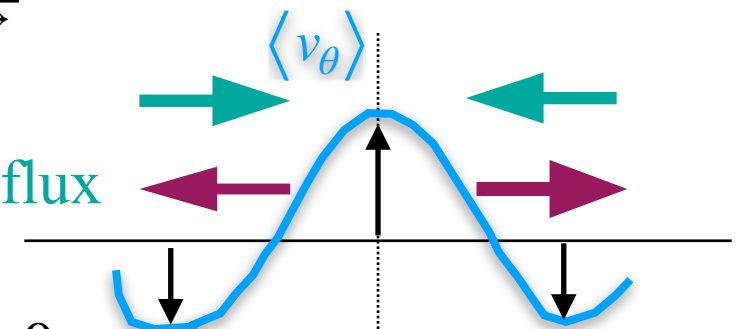
- Drift wave instability is localized at mode rational surfaces $\vec{k} \cdot \vec{B} = 0$.



- DW frequency $\omega = \frac{v_{*e} k_y}{1 + k_{\perp}^2 \rho_s^2}$, $\implies v_{gr} = -\frac{2\rho_s^2 v_{*e} k_x k_y}{(1 + k_{\perp}^2 \rho_s^2)^2}$, $v_{*e} < 0 \implies v_{gr} > 0$ for $k_x k_y > 0 \rightarrow$ outgoing wave energy

- Momentum flux: $\langle v_{E,r} v_{E,\theta} \rangle = -\frac{c}{B^2} \sum_{\vec{k}} k_x k_y |\phi_k|^2 = -\sum_{\vec{k}} v_{gx,k} \mathcal{P}_k < 0$

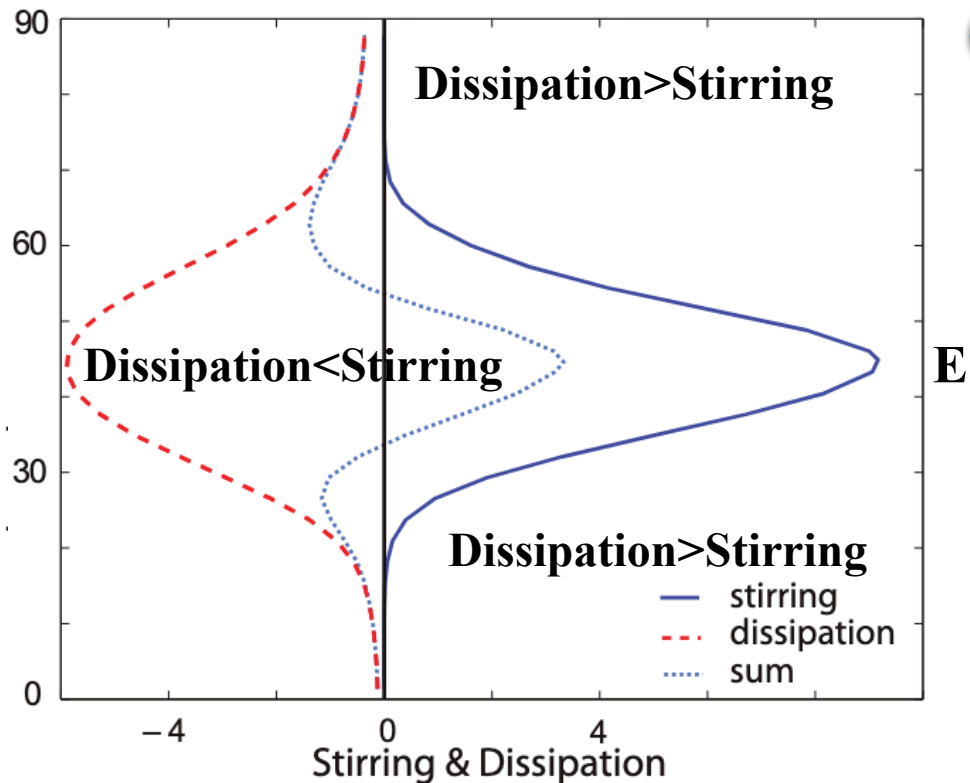
- Outgoing wave energy flux \rightarrow incoming wave momentum flux



- Zonal Flow spins up at the DW excitation region $\vec{k} \cdot \vec{B} = 0$

What determines the zonal flow pattern?

GFD Perspective



- Meridionally confined stirring.
- dissipation broader than that of forcing.
- stirring + dissipation \rightarrow zonal wind (east-ward) in the region of the stirring.

Zonal flow scale \sim wave- turbulence
 cross over scale \sim Rhines scale
 $\vec{v} \cdot \vec{\nabla} \zeta : \beta v_y \rightarrow l_R = \left(\frac{U}{\beta}\right)^{1/2}$

1. Vorticity: $\partial_t \zeta + \vec{v} \cdot \vec{\nabla} \zeta + \beta v_y = F_\zeta - D_\zeta$

2. Momentum theorem: $\partial_t \left(\mathcal{P} + \langle v_x \rangle \right) = \gamma^{-1} \left[\langle \tilde{\zeta} \tilde{F}_\zeta \rangle - \langle \tilde{\zeta} \tilde{D}_\zeta \rangle - \partial_y \langle \tilde{v}_y \tilde{\zeta}^2 \rangle \right] - r \langle v_x \rangle$

↑
↑
↑
↑
↑

Pseudomomentum
Drive
Dissipation
PE flux/
Turbulence spreading
Drag

Pseudomomentum

$$\mathcal{P} = \langle \tilde{\zeta}^2 \rangle / 2\gamma$$

$$\gamma = \beta - \partial_{yy} \langle v_x \rangle$$

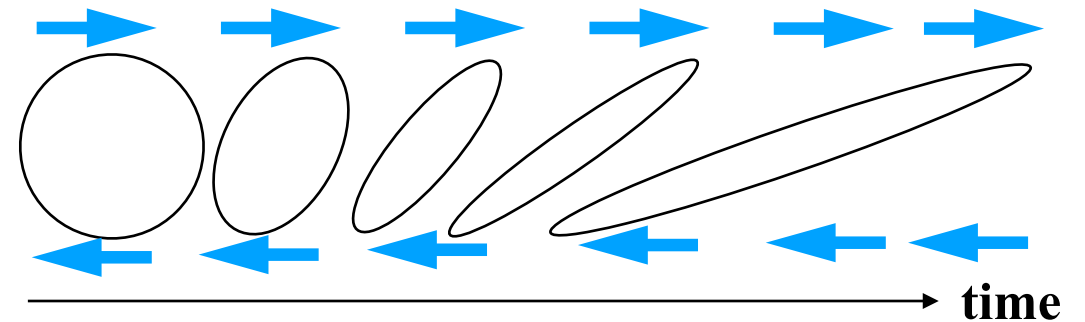
- **Non-acceleration theorem:** Absent driving (flux) and local PE decrement \rightarrow can not accelerate/maintain ZF with stationary fluctuations!

Shearing effects

- Coherent shearing [Taylor, Dupree, BDT 90]

- Radial scattering + mean shear $V'_E \rightarrow$ hybrid decorrelation

$$k_r D_{\perp}^2 \rightarrow (k_{\theta}^2 V_E'^2 D_{\perp} / 3)^{1/3} = 1/\tau_c$$



- Wave kinetics:

- $$\frac{\partial k_r}{\partial t} = - \frac{\partial (\omega + k_{\theta} V_E)}{\partial r}$$

- Mean shearing: $k_r = k_{r0} - k_{\theta} V_E' \tau_c$

- Random shearing:

➔ induced diffusion $\langle \delta k_r^2 \rangle = D_{kk} \tau$,
$$D_{kk} = \sum_q k_{\theta}^2 \left| \tilde{V}'_{E,q} \right|^2 \tau_{k,q}$$

- Mean field WKE
$$\frac{\partial \langle N_k \rangle}{\partial t} = \frac{\partial}{\partial k_x} \left[D_{kk} \frac{\partial \langle N_k \rangle}{\partial k_x} \right] + \langle \gamma_k N_k \rangle - \langle C(N_k) \rangle$$

N_k : Wave action density

Shearing effects

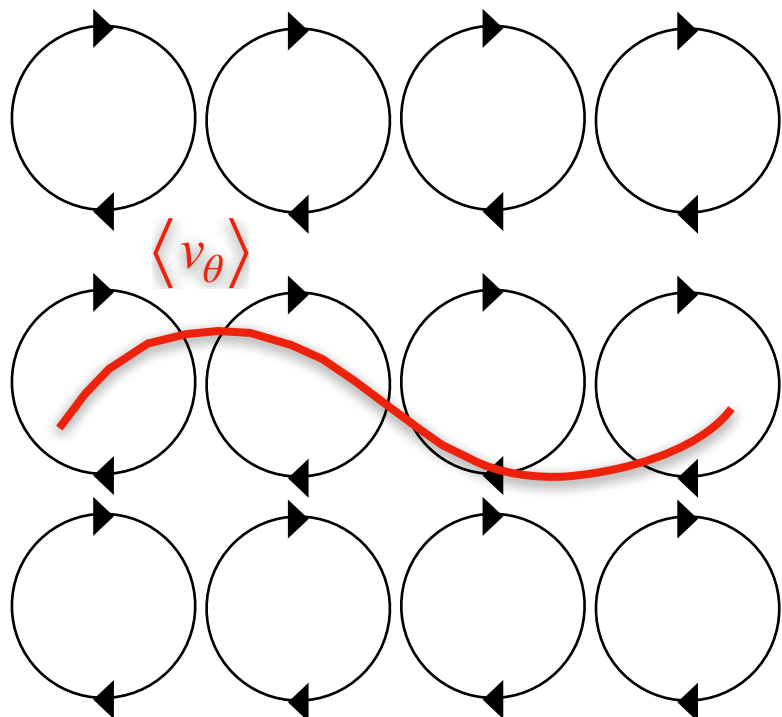
- Zonal flow shearing damps wave energy $\mathcal{E} = \int d\vec{k} \omega \langle N \rangle : \frac{\partial}{\partial t} \mathcal{E} = - \int d\vec{k} v_{gr,k} D_{kk} \frac{\partial}{\partial k_r} \langle N \rangle$

- Reduction of turbulence energy appears as growth of zonal flow

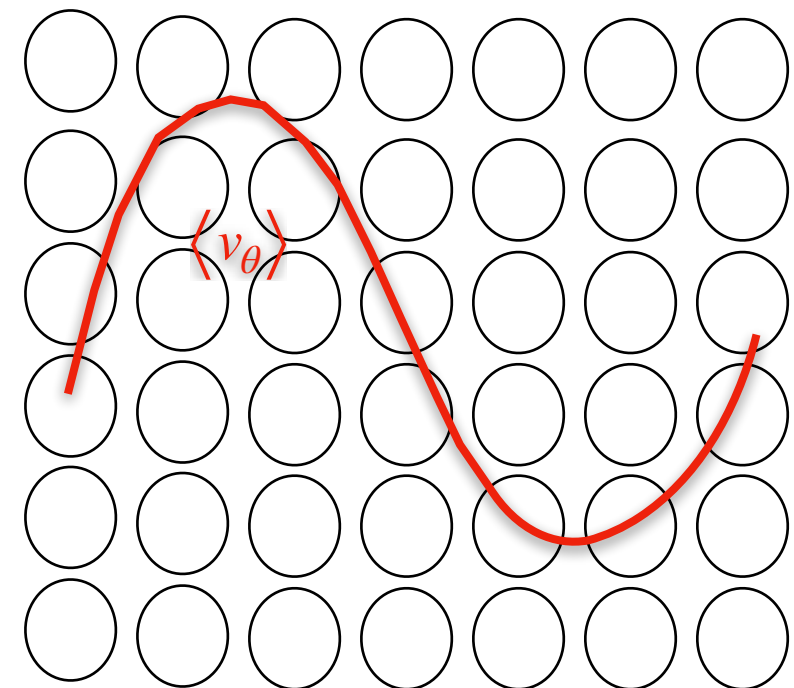
$$\frac{\partial}{\partial t} \delta \langle v_\theta \rangle + \frac{\partial}{\partial r} \delta \langle v_r v_\theta \rangle = -\nu_c \delta \langle v_\theta \rangle$$

$$\delta \langle v_r v_\theta \rangle = \int d\vec{k} k_r k_\theta \delta |\phi_k|^2 = \int d\vec{k} k_r k_\theta C \delta N_{k,q} = \int d\vec{k} k_y^2 C_k \mathcal{R}_{k,q}^{(r)} k_x \frac{\partial \langle N_k \rangle}{\partial k_x} \frac{\partial \delta \langle v_\theta \rangle}{\partial r}$$

- Modulational growth due to **-ve turbulent viscosity** for $k_x \frac{\partial \langle N_k \rangle}{\partial k_x} < 0$



Modulational growth



Noise meets modulation

- Almost all theoretical models of zonal flow generation divide cleanly into:
 1. Zonal flow dielectric or screening response, with occasional mention of **polarization beat noise**. [RH 1998, HR, 1999].

$$\frac{\partial}{\partial t} \left\langle |\phi_q|^2 \right\rangle = \frac{2\tau_c \left\langle |S_q|^2 \right\rangle}{\chi(q)} \quad S_q = \text{Polarization flux}$$

$$\chi = \chi_{cl} + \chi_{neo} = \left\{ 1 + \frac{q^2}{\epsilon^2} \right\} k_r^2 \rho_i^2 \quad ;$$

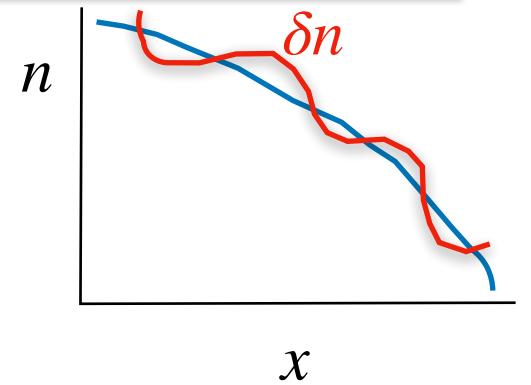
-overlooks modulational mechanism - "negative viscosity"

2. Modulational stability calculations ignore noise emission.
- Interaction? → Coupled spectral evolution - DIA closure
 - What of profile corrugations ?
 - What of zonal shears and corrugations alignment - staircases?
 - Noise effect on feedback processes ?
 - Macroscopics - L→H transition?

Profile corrugations meet zonal flows

- Wave kinetics:

- $\partial_t k_r = -\partial_r (\omega + k_\theta V_E); \quad \omega = \omega(n', T', \dots)$



- Freq. modulation due to **profile corrugations** $\delta\omega = \frac{\partial\omega}{\partial n'}\delta n'' + \frac{\partial\omega}{\partial T'}\delta T'' + \dots$

- Random shearing: \rightarrow induced diffusion $\langle \delta k_r^2 \rangle = D_{kk}\tau$,

$$D_{kk} = \sum_q k_\theta^2 \left| \tilde{V}'_{E,q} + \frac{\partial\omega}{\partial n'}\delta n'' + \frac{\partial\omega}{\partial T'}\delta T'' + \dots \right|^2 \tau_{k,q}$$

➡ Induced diffusion depends on the relative alignment of zonal shear and profile corrugations!

➡ Interferes fluctuation energy and zonal flow saturation !

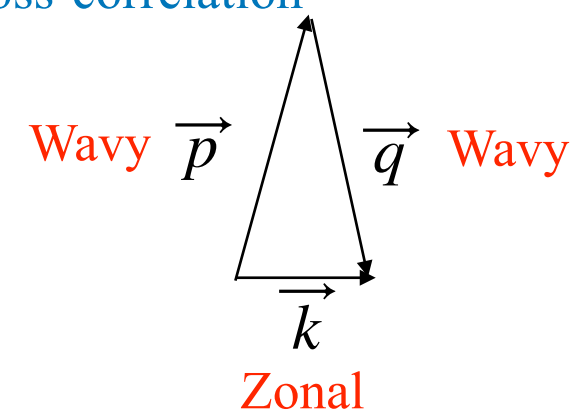
➡ ... unified theory of zonal flow and corrugations needed! Use Hasegawa- Wakatani model at the simplest level of description.

Spectral evolution of zonal intensity

For the zonal mode $k_y = k_{\parallel} = 0$ and $k_x \neq 0$

$$\left(\frac{\partial}{\partial t} + 2\mu k_x^2\right) \langle |\phi_k|^2 \rangle + 2\eta_{1k}^{\text{zonal}} \langle |\phi_k|^2 \rangle + \Re \left[2\eta_{2k}^{\text{zonal}} \langle n_k \phi_k^* \rangle \right] = F_{\phi k}^{\text{zonal}}$$

Zonal density - potential cross-correlation



- NL damping rate: $\eta_{1k}^{\text{zonal}} \propto k_x^2$ and becomes -ve for $\frac{\partial I_q}{\partial q_x} < 0 \rightarrow$ transfer to large scales by **negative turbulent viscosity!**

- Zonal growth is maximum when the adiabaticity parameter $\alpha_q \rightarrow \infty \Rightarrow$ **Non-adiabatic fluctuations inhibit transfer to large scales !**

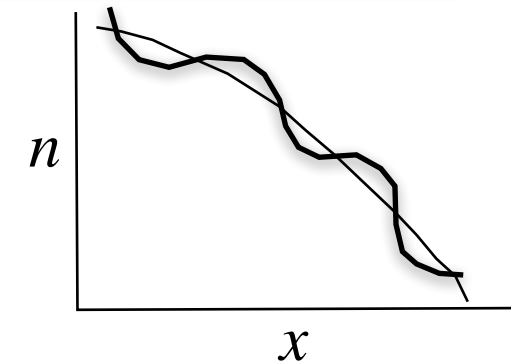
- Cross transfer rate: $\eta_{2k}^{\text{zonal},(r)} > 0$ ALWAYS for $\frac{\partial I_q}{\partial q_x} < 0 \Rightarrow$ Forward transfer when $\Re \langle n_k \phi_k^* \rangle < 0$, backward transfer when $\Re \langle n_k \phi_k^* \rangle > 0$

- Noise: **always +ve and of envelope scale!** $F_{\phi k}^{\text{zonal}} = 4 \sum_q \Pi_q^2 \Theta_{k,-q,q}^{(r)}$; Reynolds stress $\Pi_q = q_y q_x I_q$. Noise/Modulation = $q_x^2 I_q / k_x^2 I_k =$ Turbulent KE/ Zonal KE.

Spectral evolution of density corrugations

[Singh, Diamond PPCF 2021]

$$\left(\frac{\partial}{\partial t} + 2D_n k^2\right) \langle |n_k|^2 \rangle + 2\zeta_{1k} \langle |n_k|^2 \rangle + \Re [2\zeta_{2k} \langle n_k^* \phi_k \rangle] = F_{nk}$$



Corrugations
damping rate,
+ve $\sim 1/\alpha_q^2$

Cross transfer
rate, +ve
 $\sim 1/\alpha_q^2$

Advection
noise, +ve
 $\sim 1/\alpha_q^2$

- Density cascade forward in k_x !
- Corrugations become weaker as the response become more adiabatic.
- Corrugation is determined by noise vs diffusion balance.
- Important for the nonlinear dynamics underlying staircases. Forward cascade in k_x -space is supporting the idea of (inhomogeneous) mixing in real space.

Zonal cross-correlation(ZCC)

[Singh, Diamond PPCF 2021]

- Significant for layering or staircase structure - ZF and ∇T are aligned in staircase!

- When do zonal density and zonal potential align?

$$\Re \langle n_k \phi_k^* \rangle = \frac{2\eta_{2k}^{(r)} \langle |n_k|^2 \rangle + 2\zeta_{2k}^{(r)} \langle |\phi_k|^2 \rangle}{-(\mu + D_n) k_x^2 - 2\xi_{1k}^{(r)}}; \quad \xi_{1k}^{(r)} = \eta_{1k} + \zeta_{1k}$$

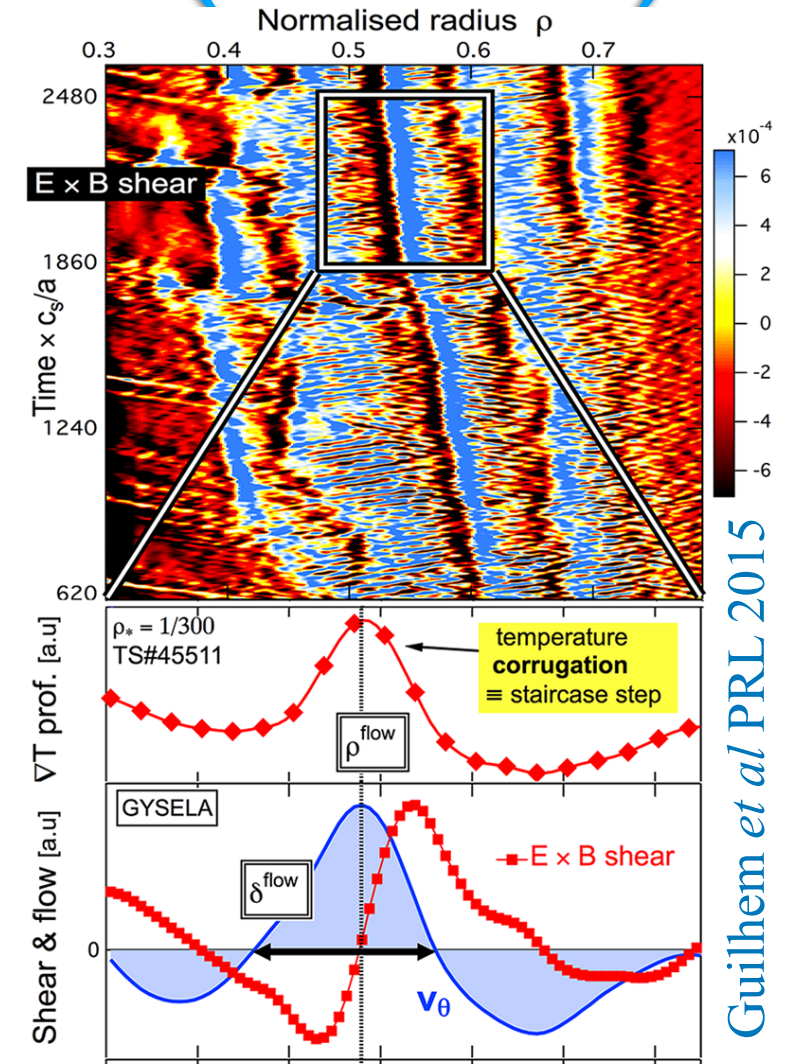
- Zonal density and potential are correlated (anti-correlated) when the modulational growth of zonal flow more (less) than modulational damping of corrugations.

- Imposing physically bounded solution for $\Im \langle n_k \phi_k^* \rangle = 0$ fixes the sign of $(\mu + D_n) k_x^2 + 2\xi_{1k}^{(r)} > 0$. Thus $\Re \langle n_k \phi_k^* \rangle < 0$.

- Hence ZCCs in real space are: $\langle \bar{n} \bar{\phi} \rangle < 0$, $\langle \bar{n} \nabla_x^2 \bar{\phi} \rangle > 0$, $\langle \nabla_x \bar{n} \nabla_x^2 \bar{\phi} \rangle = 0$

- $\langle \nabla_x \bar{n} \nabla_x^3 \bar{\phi} \rangle > 0$: zonal density jumps are co-located with the zonal vorticity jumps.

- $\langle -\nabla_x \bar{n} \nabla_x \bar{\phi} \rangle > 0$: density gradient peaks are co-located with the zonal flow peaks.



Guilhem et al PRL 2015

Summary of zonal flow and corrugations interaction

(a) Zonal flow - Vorticity equation - Polarization charge flux		
Process	Impact	Key physics
Polarization noise	Seeds zonal flow	Polarization flux correlation, +ve definite
Zonal flow response (comparable to noise)	Drives zonal shear using DW energy	Non-local inverse transfer in k_x , -ve viscosity
Zonal shear straining of small scale	Regulates waves via straining	Stochastic refraction straining waves, induced diffusion to high k_x
(b) Density corrugations - Density equation - Particle flux		
Density advection beat noise	Seeds density corrugation	Advection beats due to non-adiabatic electrons.
Density corrugations response	Damps and regulates density corrugations	Non-local forward transfer in k_x +ve diffusivity, turbulent mixing weak for $\alpha \gg 1$
Zonal shear straining of small scale	Regulates waves via straining	Stochastic refraction straining waves, induced diffusion to high k_x
(c) Zonal cross-correlation - Vorticity and density transport processes		
ZCC response	Sets corrugation - shear layer correlation; staircase states	Growth of zonal intensity must exceed the modulational damping rate of corrugation

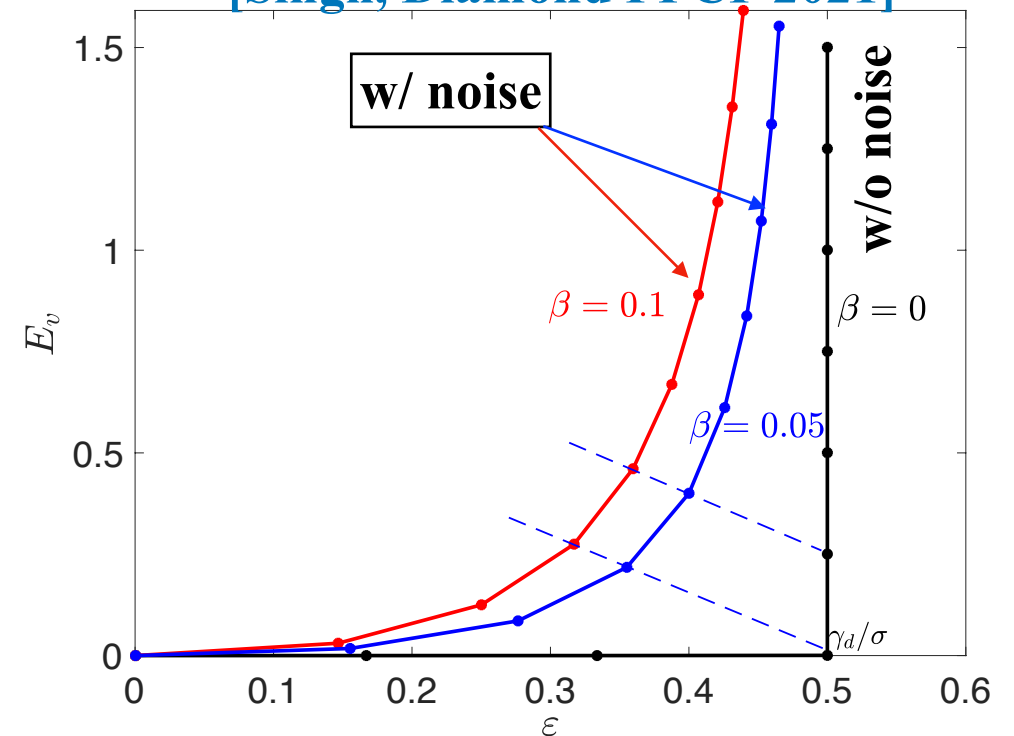
Feedback loop with nonlinear zonal noise

[Singh, Diamond PPCF 2021]

Feedback + **Noise** → revisit predator - prey

Turbulence energy \mathcal{E} : $\frac{\partial \mathcal{E}}{\partial t} = \gamma \mathcal{E} - \underbrace{\sigma E_v \mathcal{E}}_{\text{Induced diffusion}} - \underbrace{\eta \mathcal{E}^2}_{\text{Nonlinear damping}}$
(Prey)

Zonal flow energy E_v : $\frac{\partial E_v}{\partial t} = \underbrace{\sigma \mathcal{E} E_v}_{\text{Modulational growth}} - \gamma_d E_v + \beta \mathcal{E}^2$
(Predator)



Without noise:

- Threshold in growth rate $\gamma > \eta \gamma_d / \sigma$ for appearance of stable zonal flows.
- Turbulence energy increases as γ / η below the threshold, until it locks at γ_d / σ , at the threshold.
- Beyond the threshold, turbulence energy remains locked at $\frac{\gamma_d}{\sigma}$ while the zonal flow energy continues to grow as $\sigma^{-1} \eta (\gamma / \eta - \gamma_d / \sigma)$.

With noise:

- Both zonal and turbulence co-exist at any growth rate - No threshold in growth rate for zonal flow excitation.
- Turbulence energy never hits the old modulational instability threshold, absent noise!
- Turbulence energy ↓ and zonal flow energy ↑:- Noise feeds energy into zonal flow!

L - H transition

How does zonal noise affect the dynamics of L-H transition ?

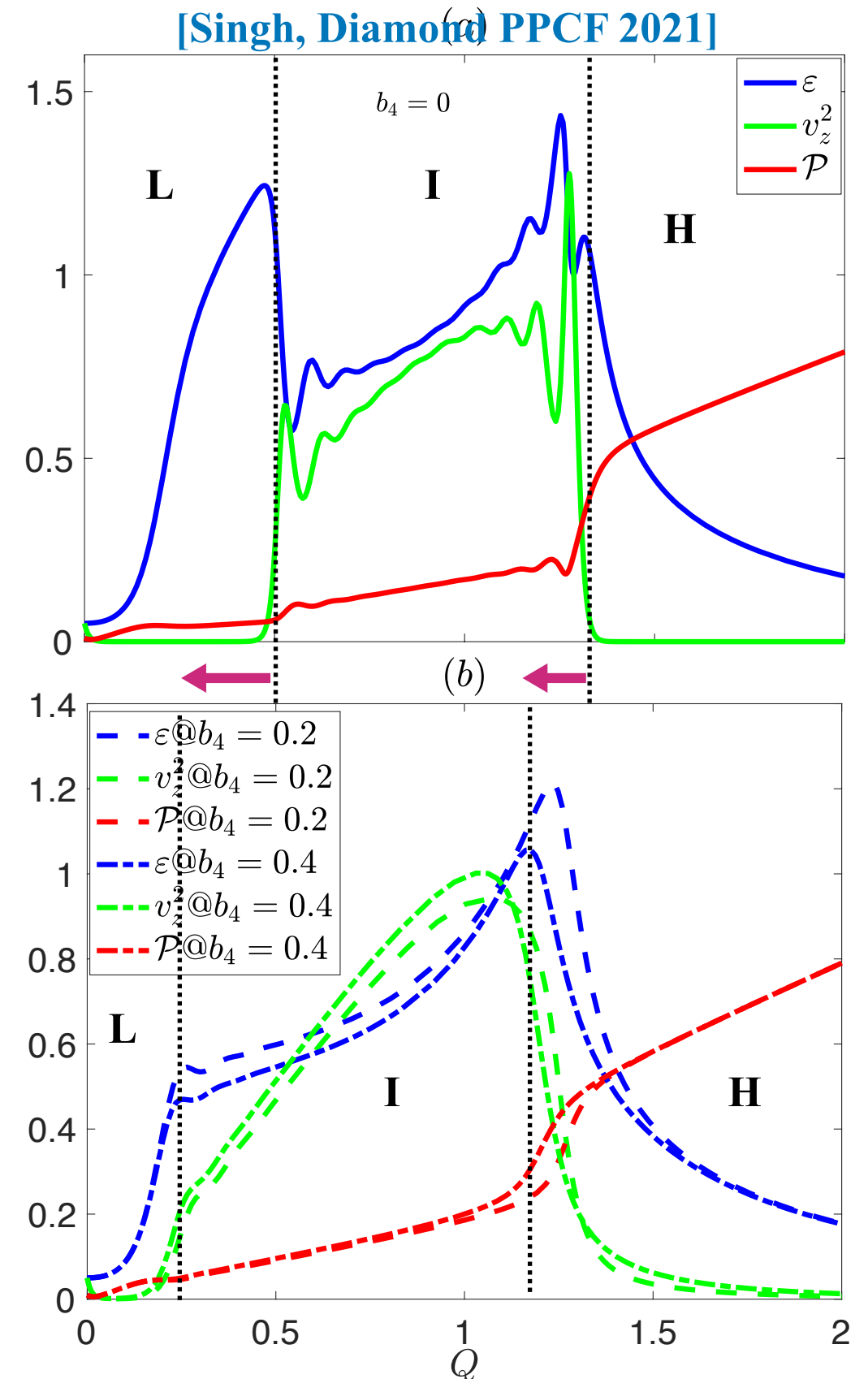
KD03 + Noise

Without noise:

- Critical power, set by modulational instability threshold.
- Zonal flows exist only within the I-phase.

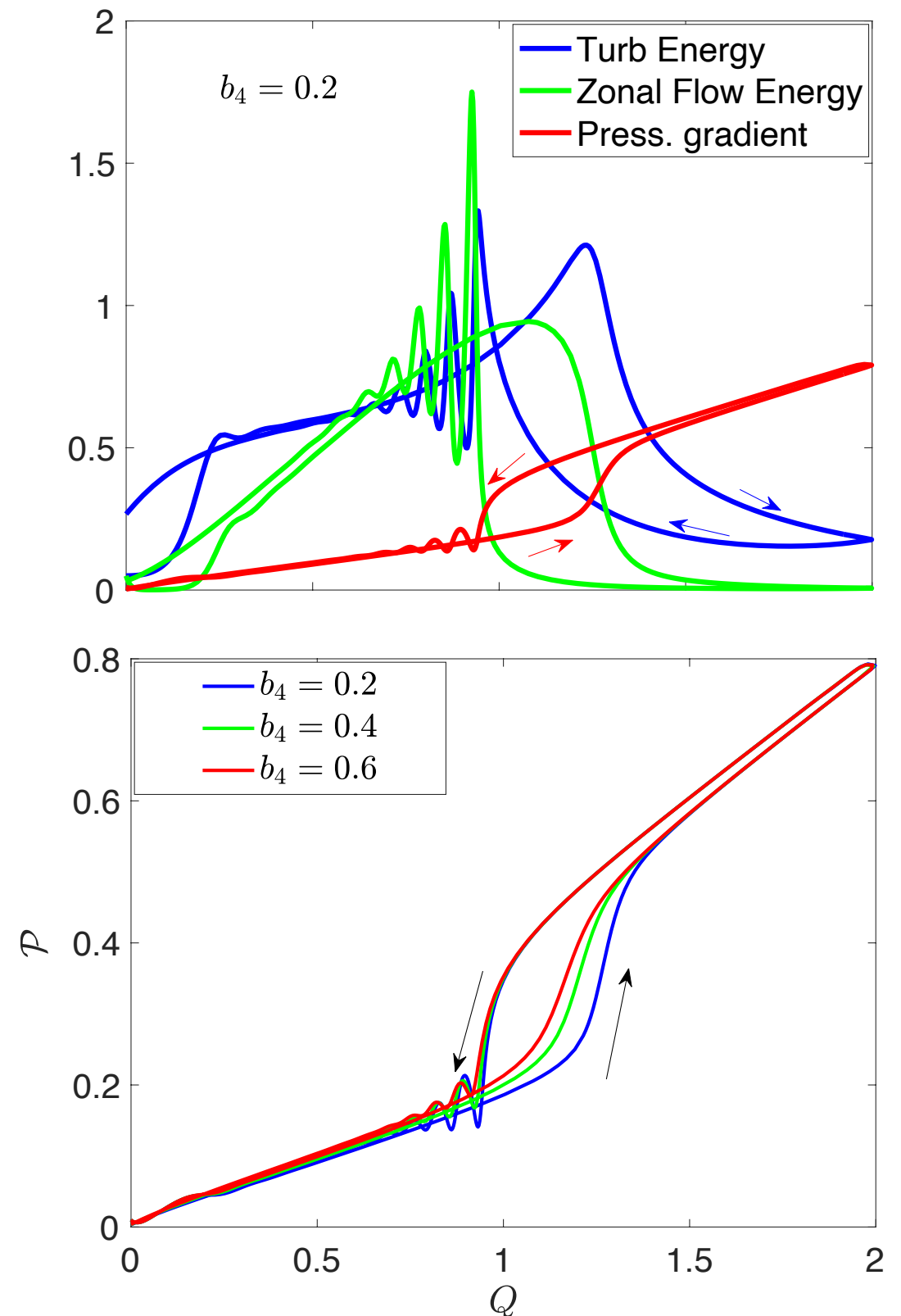
With Noise: No ZF threshold in Q !

- Turbulence level \downarrow , no overshoot, zonal flow \uparrow .
- I-Phase oscillations \downarrow .
- H-mode power threshold \downarrow . Noise couples more fluctuation energy to zonal flows.



Noise effect on L-H-L hysteresis

- A cyclic power ramp exhibits hysteresis.
- The I-phase in the back transition is more oscillatory than that in the forward transition.
- Hysteresis with noise is robust w.r.t the variations in the initial conditions.
- Threshold power for both forward and backward transition decreases with noise such that the area enclosed the hysteresis curve decreases with noise.



Conclusions I

- Unified theory of zonal modes (flows and corrugations) encompassing both noise and modulations.
 - Vorticity flux corr. \rightarrow ZF noise. density flux corr. \rightarrow corrugation noise.
 - Bi-directional transfer: KE energy to large scales, internal energy to small scales. Turbulent viscosity -ve but turbulent diffusivity +ve.
 - The effective zonal viscosity goes negative only for an energy spectrum which decays sufficiently rapidly in k_r i.e., $\partial E / \partial k_r < 0$ and $\left| \partial E / \partial k_r \right| < \left| \partial E / \partial k_r \right|_{crit}$.
 - Zonal cross-correlation $\Re \langle n_k \phi_k^* \rangle < 0 \rightarrow \langle \nabla \bar{n} \nabla^3 \bar{\phi} \rangle > 0$ i.e., zonal density gradient jumps are colocated with zonal vorticity jumps.

Conclusions II

Implications:

- Polarization beat noise and modulational effects are comparable intrinsically (both driven by Reynolds stress!). The synergy of the two mechanisms is stronger than either alone.
 - Expands the range of zonal flow activity relative to that predicted by modulational instability calculations.
 - Increases branching ratio of zonal flow energy to turbulence energy.
- Interaction of zonal noise and modulation: → significant effect on feedback processes !
 - L-H transition: Noise eliminates the threshold for zonal flow excitation, and so expands the predicted range of the intermediate phase, drastically reduces the turbulence overshoot.
 - The energy transfer to zonal flow is accelerated which lowers the threshold for L-H transition.

Future directions

- Understanding interaction of corrugations with avalanches:
 - Corrugations in state of high ZCC sustained as localized transport barriers, staircases etc. localized by accompanying shear flow?
 - Corrugations in state of low ZCC likely to overturn, and drive avalanches, as in running sandpile?
 - Relevant for TEM turbulence. Does the density gradient state consist of standing corrugations , running avalanches or mixtures thereof ?
- Effect of noise on staircase? (have been considered only in context of Mean Field theory).
- Relation between ZCC and the staircase structure: Does the physics of ZCC set the relative positions of corrugations and shear layer? Is there a single ZCC for staircase state ? Or a band ?
- GFD: Are zonal wind and temperature profile correlated? ZCC and weather pattern?